



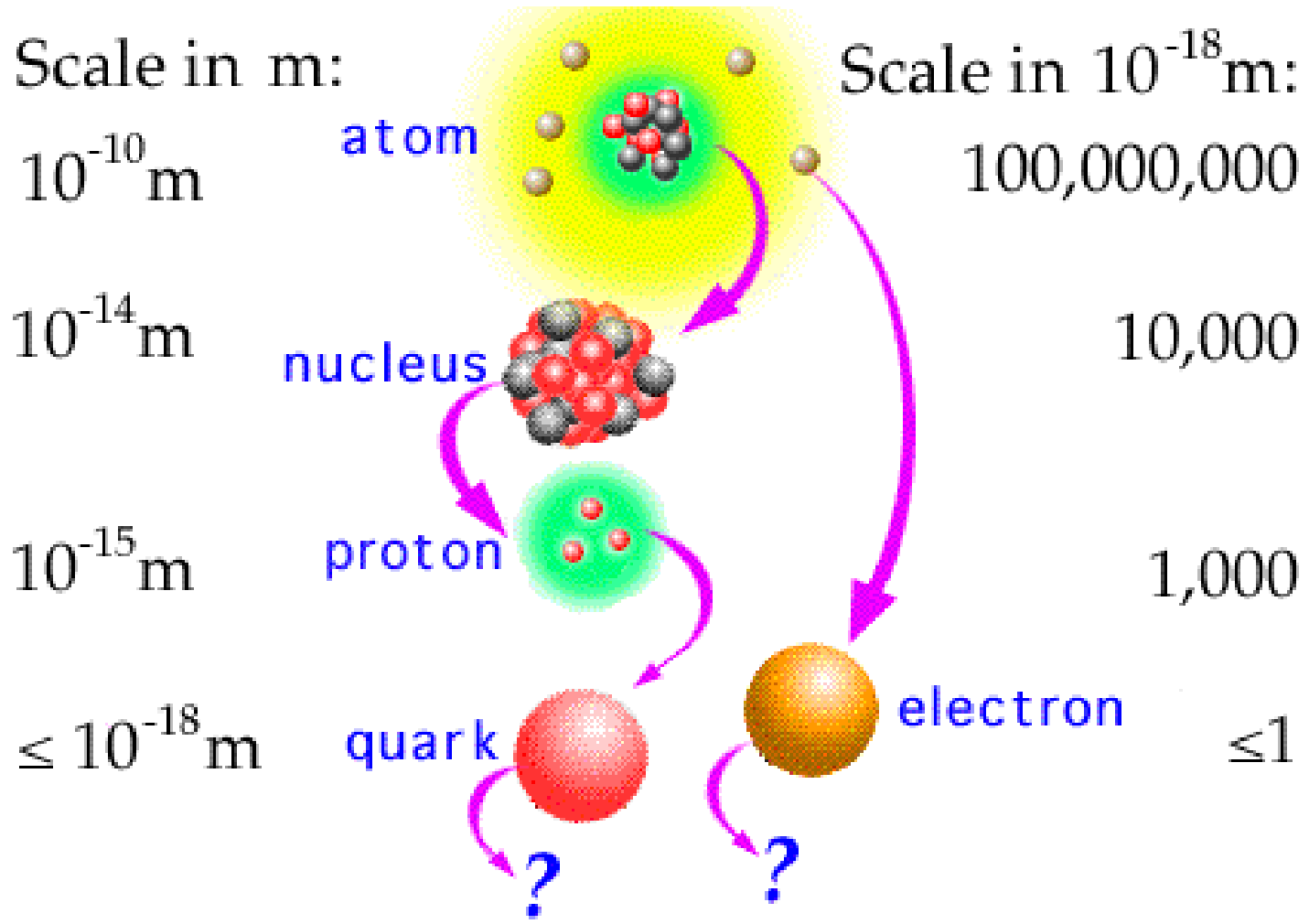
Nuclear Interactions

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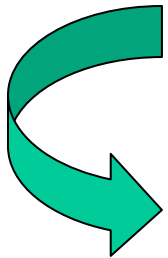
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Scales in Nature



Energy Densities:

- Spring under tension
- Binding energy of electrons in a heavy atom
- Binding energy of nucleons
- $\approx 10^6 \text{ J/m}^3$
- $\approx 10^{16} \text{ J/m}^3$ [eV]
- $\approx 10^{32} \text{ J/m}^3$ [MeV]



Nuclear force → strong interaction

Travel through History

1930's:

- Chadwick (1932) discovers the neutron
- Heisenberg (1932) proposes the first phenomenology (Isospin)
- Yukawa (1935) proposes the pion and his Meson Hypothesis

1940's:

- Discovery of the pion in cosmic rays (1947)
- and in the Berkeley Cyclotron Lab (1948)
- Nobel Prize awarded to Yukawa (1949)
- Rabi (1948) measures the quadrupole moment of the deuteron

Low Energy Nuclear Physics

Smallest Nucleus: Deuteron (np)

- Ground State

Properties

- Binding Energy
- Spin & Parity
- Iso Spin (n,p)
- Magnetic Dipole Moment
- Electric Quadrupole Moment
- Radius

- Value

2.22457312(22) MeV

1+

0

0.857406(1) μ_N

0.28590(3) fm²

1.963(4) fm

Nuclear Force has Tensor Component

- Yukawa 1934:
 - postulated π as carrier of the nuclear force

$$-\frac{\vec{\sigma}_1 \cdot \vec{p} \vec{\sigma}_2 \cdot \vec{p}}{m_\pi^2 + p^2} \approx \left\{ \frac{3}{r^2} (\vec{\sigma}_1 \cdot \vec{r} \vec{\sigma}_2 \cdot \vec{r}) - (\vec{\sigma}_1 \cdot \vec{\sigma}_2) \right\} \frac{e^{-m_\pi r}}{r}$$

- Analogy EM: Quadrupole momentum tensor:

$$Q_{ij} = \int \rho(x) (3x_i x_j - r^2 \delta_{ij}) d^3x ,$$

- Consequence: $[H, L_3] \neq 0$ but $[H, J_3] = 0$

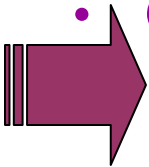
More experimental evidence

- Comparison of the binding energies of:
 - 2H (deuteron), 3H (triton), 4He (alpha - particle)
 - show that the nuclear force is of finite range (1 – 2 fm) and very strong within that range.
- For nuclei with $A > 4$:
 - the energy saturates: Volume and binding energies of nuclei are proportional to the mass number A .
 - Binding energy of all nuclei $\sim A^{1/3}$
 - Nuclear force must be repulsive at short distances
- Spin-orbit force:
 - Observation of large polarizations of scattered nucleons perpendicular to the plane of scattering.

Travel through History

1950's:

- Taketani, Nakamura, Sasaki (1951):
 - One Pion Exchange (OPE) o.k.
 - Multi-Pion Exchanges – problems!
- Taketani, Machida, Onuma (1952): multi-pion exchanges
 - Concept of renormalization was not worked out
- Gammel-Thaler (1957): First NN potential model
- Okubo-Marshak (1958) [Ann.Phys. 4, 166]
 - General consideration about the form of an NN interaction using only invariance properties of a Hamiltonian.



Travel through History

1960's:

- Many pions = multi-pion resonances:
 - $\sigma(600)$, $\rho(770)$, $\omega(782)$
 - Input from High Energy Physics
- One-Boson-Exchange Models
 - General Invariance structure given by Poincaré Invariance.
- Refined phenomenology: Reid Potential (1968)

1970's and 1980's:

- Sophisticated models for two-pion exchange
 - Paris Potential (Lacombe et al., Phys. Rev. C 21, 861 (1980))
 - Bonn Potential (Machleidt, Holinde, Elster, Phys. Rep. 149, 1 (1987))
- Quark Cluster Models
- Begin of Effective Field Theory studies

Travel through History

1990's:

- “High-Precision” NN Potentials:
 - Nijmegen I, II, '93, Reid93 (Stoks et al. 1994)
 - Argonne V18 (Wiringa et al, 1995)
 - CD-Bonn (Machleidt et al. 1996 and 2001)
- Advances in Effective Field Theory – Chiral Perturbation Theory
 - Weinberg (1990)
 - Ordonez, Ray, van Kolck (1996) ...

2000 + :

- Multi-Pion theory constrained by Chiral Symmetry:
 - Chiral EFT (lots of ongoing work)
- 2006: Nucleon-nucleon interaction from Lattice QCD,
final confirmation of meson hypothesis of Yukawa?

Okubo – Marshak Invariants

NN Hamiltonian:
$$H = \frac{1}{2m} (p_1^2 + p_2^2) + V_{NN}(1,2)$$

$$V_{NN}(1,2) \equiv V_{NN}(\vec{r}, \vec{R}, \vec{p}, \vec{P}, \vec{\sigma}^{(1)}, \vec{\sigma}^{(2)}, \vec{\tau}^{(1)}, \vec{\tau}^{(2)}) \quad (V_{NN} \text{ Hermitian})$$

- Translation invariance
- Rotation invariance
- Space reflection invariance
- Time reversal invariance
- Galilean invariance
- Invariance under the exchange of nucleons 1 and 2
- Almost Isospin Symmetry

(www.phy.ohiou.edu/~elster/lectures/fewblect_1.pdf Section 1.9)



Most general, velocity-dependent, non-relativistic NN potential

$$V_{NN} \equiv V_{NN}(\vec{r}, \vec{R}, \vec{p}, \vec{P}, \vec{\sigma}^{(1)}, \vec{\sigma}^{(2)}, [\vec{\tau}^{(1)} \cdot \vec{\tau}^{(2)}])$$

$$= V_1(\vec{r}, \vec{R}, \vec{p}, \vec{P}, \vec{\sigma}^{(1)}, \vec{\sigma}^{(2)}) + (\vec{\tau}^{(1)} \cdot \vec{\tau}^{(2)}) V_2(\vec{r}, \vec{R}, \vec{p}, \vec{P}, \vec{\sigma}^{(1)}, \vec{\sigma}^{(2)})$$

with

$$\begin{aligned} V_i &= V_i^c(r^2, p^2, L^2) + \vec{\sigma}^{(1)} \cdot \vec{\sigma}^{(2)} V_i^\sigma(r^2, p^2, L^2) \\ &+ S_{12} V_i^T(r^2, p^2, L^2) + \vec{S} \cdot \vec{L} V_i^{LS}(r^2, p^2, L^2) \\ &+ [(\vec{\sigma}^{(1)} \cdot \vec{L})(\vec{\sigma}^{(2)} \cdot \vec{L}) + (\vec{\sigma}^{(2)} \cdot \vec{L})(\vec{\sigma}^{(1)} \cdot \vec{L})] V_i^{\sigma L}(r^2, p^2, L^2) \end{aligned}$$

and
$$S_{12} = \frac{3}{r^2} (\vec{\sigma}^{(1)} \cdot \vec{r})(\vec{\sigma}^{(2)} \cdot \vec{r}) - (\vec{\sigma}^{(1)} \cdot \vec{\sigma}^{(2)})$$

e.g. one-pion-exchange (OPE):

$$V_{\text{OPE}}(\mathbf{r}) \equiv \frac{1}{12} \frac{g_{ps}^2}{4\pi} m_{ps} \left[\left(\frac{m_{ps}}{m} \right)^2 \mathcal{Y}(m_{ps}r) \vec{\sigma}_1 \cdot \vec{\sigma}_2 + Z(m_{ps}r) S_{12} \right]$$

$$\mathcal{Y}(x) = e^{-x}/x$$

$$Z(x) = (m_{\alpha}/m)^2 (1 + 3/x + 3/x^2) \mathcal{Y}(x)$$

Momentum space:

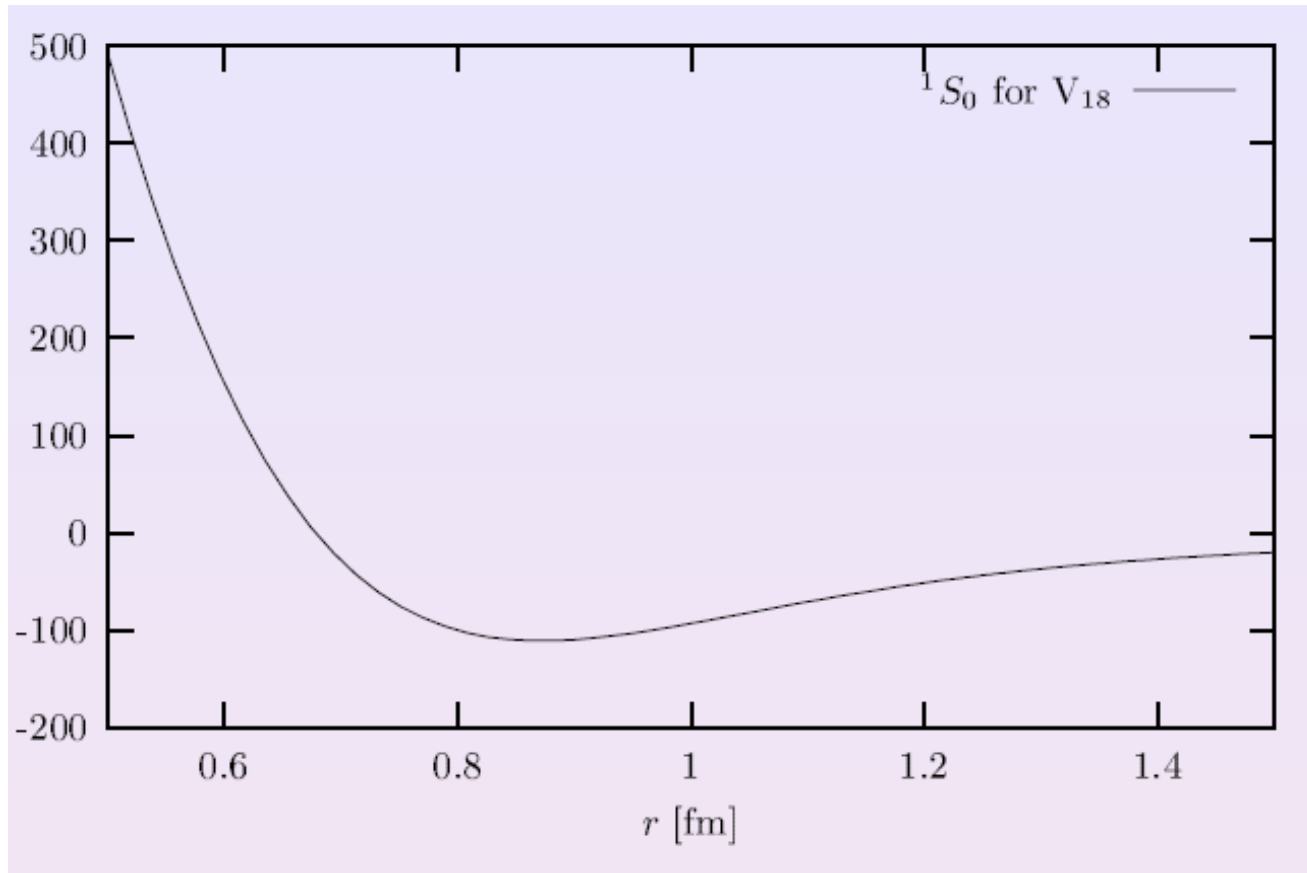
$$V_{\text{OPE}}(\mathbf{k}) = - \frac{g_{ps}^2}{4m^2} \frac{(\vec{\sigma}_1 \cdot \vec{k})(\vec{\sigma}_2 \cdot \vec{k})}{k^2 + m_{ps}^2}$$

Task:

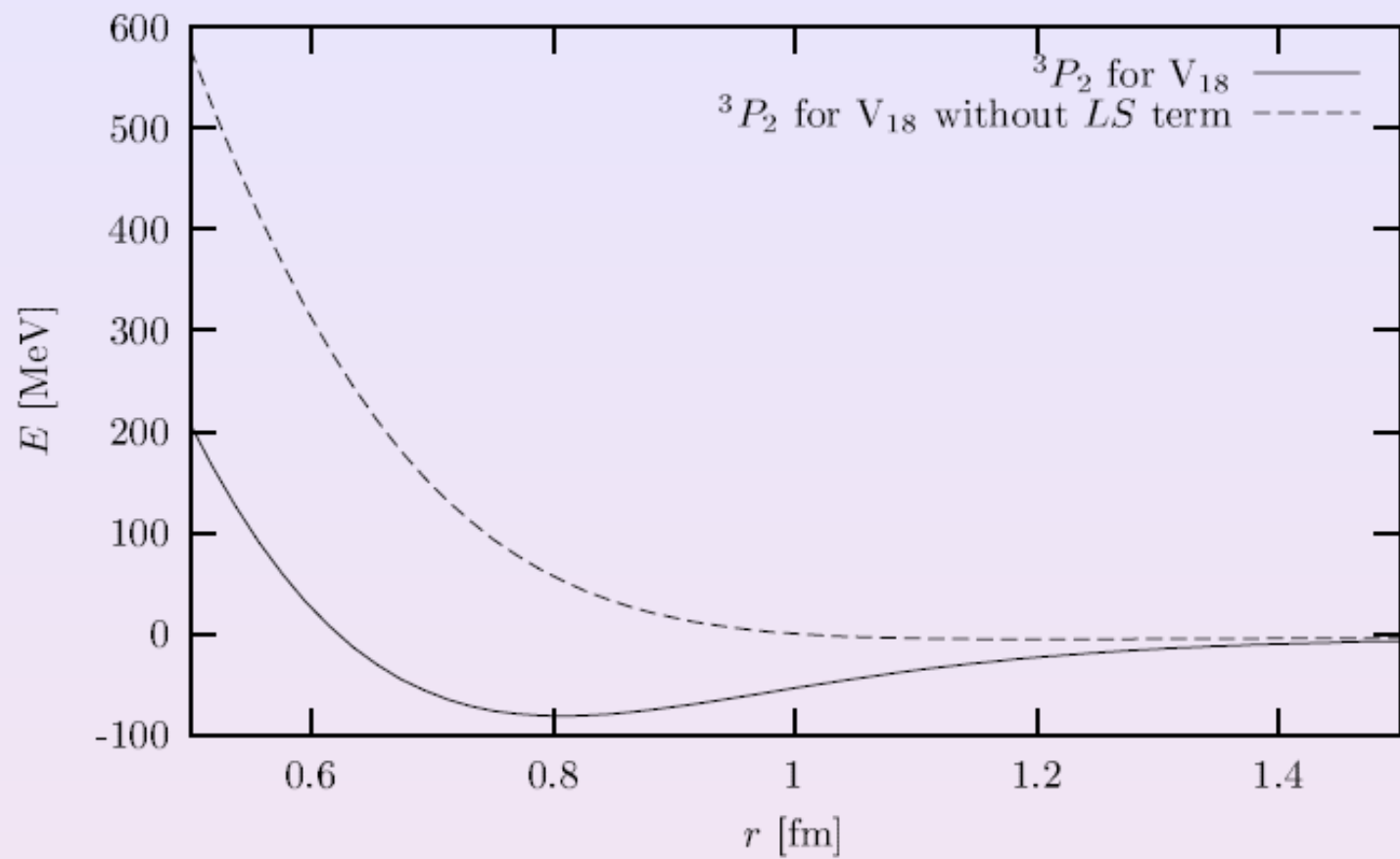
Derive / obtain functions $V_i^\alpha (r^2, p^2, L^2)$:

- **Ansatz I** (based on Okubo-Marshak Invariants)
 - Functions: Yukawa type $\sim \text{const} * \exp(-m_\alpha r) / r$
 - $\text{const} \equiv$ magnitude of term; $m_\alpha \equiv$ range of term
 - For longest range: OPE
 - Choose superposition of terms with different strengths and ranges and fit to NN data
 - NN data:
 - cross section, polarization and double polarizations
 - Potentials:
 - Reid (1968)
 - Nijmegen potentials
 - Argonne V18

Potential projected on angular momentum eigenstates




Notation: $^{2S+1}L_J$



Ansatz II

- Start from relativistic Lagrangians
 - Invariants: bilinear covariants (5 quantities invariant under Lorentz group)
- Input: mesons from the particle data (mesons $\equiv \langle q\bar{q} \rangle$ pairs)
 - Pseudo-scalar / scalar / vector mesons
- Functional forms of the $V_i^\alpha (r^2, p^2, L^2)$ given
 - From evaluation of Feynman diagrams with exchanged meson
- Parameters
 - Coupling constants, cutoff parameter(s)
- One-boson-exchange (OBE) potentials
 - Bonn series (A,B,C)
 - Utrecht models
 - Gross model
 - CD-Bonn (2001)

Dramatis Personae:

Mesons	Mass (MeV)
π	138.03
η	548.8
 σ	≈ 550.0
ρ	770
ω	782.6
δ	983.0
K	495.8
K^*	895.0

Phenomenological Lagrangians
for spin 1/2 baryons:

$$\mathcal{L}_{ps} = g^{ps} \bar{\Psi} \gamma^5 \Psi \phi^{(ps)},$$

$$\mathcal{L}_s = g^s \bar{\Psi} \Psi \phi^{(s)},$$

$$\mathcal{L}_{pV} = \frac{f_\pi}{m_\pi} \bar{\psi} \gamma_5 \gamma_\mu \psi \partial^\mu \phi_\pi.$$

$$\mathcal{L}_V = g^V \bar{\Psi} \gamma_\mu \Psi \phi_\mu^{(V)} + g^T \bar{\Psi} \sigma^{\mu\nu} \Psi \left(\partial_\mu \phi_\nu^{(V)} - \partial_\nu \phi_\mu^{(V)} \right)$$

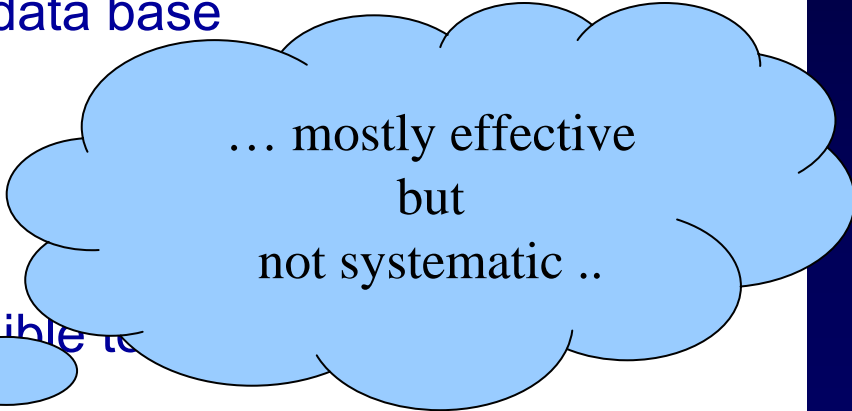
Ansatz II cont'd

- Scalar meson σ responsible for attractive piece in NN force
 - Existence questionable in the history of PDG
- Explicit 2-pion exchange into OBE models
 - Via Dispersion relations
 - Paris Potential
 - Via explicit incorporation of explicit two-meson exchanges ($\pi\pi$, $\pi\rho$, $\rho\rho$) and the Delta resonance
 - Bonn Potential (1987)

Assessment:

Potentials based on Galilei (or Poincaré) invariance and meson exchange:

- Very successful
 - Perfect description of world NN data base
- However:
 - Little relation to QCD
 - Model dependence
 - Theoretical uncertainties impossible to



... mostly effective
but
not systematic ..

Ansatz III : Chiral Effective Field Theory

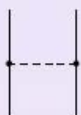
- Write down the most general Lagrangian including all terms consistent with the assumed symmetries, in particular the spontaneously broken chiral symmetry
- Calculate all Feynman diagrams
 - There will be infinitely many
- Find a scheme for assessing the importance of the various diagrams
- Organize the contributions in terms of $\Lambda_\chi \approx 1 \text{ GeV}$

$$\mathcal{L}_{NN} = \mathcal{L}_{NN}^{(0)} + \mathcal{L}_{NN}^{(2)} + \mathcal{L}_{NN}^{(4)} + \dots,$$

where the superscript refers to the number of derivatives or pion mass insertions (chiral dimension). Good review: Epelbaum, Prog. Part. Nucl. Phys. **57**, 654 (2006).

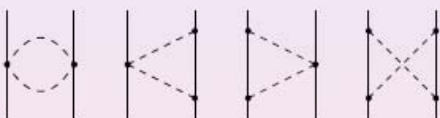
Chiral EFT

Q^0



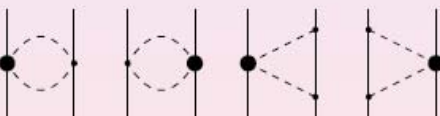
1π Exchange

Q^2



2π Exchange

Q^3



χ PT

- The most important irreducible one- and two-pion exchange contributions to the NN interaction up to order Q^3 .
- Vertices denoted by small dots are from $\widehat{\mathcal{L}}_{\pi N}^{(1)}$.
- Large dots refer to $\widehat{\mathcal{L}}_{\pi N, ct}^{(2)}$

Major advantage: Systematic expansion **and** error estimates in each order